1 GOES 8 aerosol optical thickness assimilation in a mesoscale 2 model: Online integration of aerosol radiative effects

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6 [1] To investigate the importance of aerosol radiative effects in the troposphere, numerical 7 simulation of a dust event during the Puerto Rico Dust Experiment is presented by using

8 the Colorado State University Regional Atmospheric Modeling System (RAMS).

9 Through assimilation of geostationary satellite-derived aerosol optical thickness (AOT)

into the RAMS, spatial and temporal aerosol distribution is optimally characterized,
 facilitating direct comparison with surface observations of downwelling radiative energy

12 fluxes and 2 m air temperature that is not possible with a free-running mesoscale

model. Two simulations with and without consideration of aerosol radiative effects are

14 performed. Comparisons against observations show that direct online integration of

¹⁵ aerosol radiative effects produces realistic downwelling shortwave and longwave fluxes at

the surface but minimal improvement on 2 m air temperature at the observation location.

17 Numerical simulations show that for the dust loading considered in this study (AOT =

¹⁸ 0.45 at 0.67 μ m), if the dust radiative effects are not properly represented, the uncertainty ¹⁹ in the simulated AOT is about ±5 to ±10%, the surface radiative energy is overestimated

in the simulated AOT is about ± 5 to $\pm 10\%$, the surface radiative energy is overestimated by 30-40 W m⁻² during the day and underestimated by 10 W m⁻² during the night,

and the bias in air temperatures near the surface could be up to $\pm 0.5^{\circ}$ C, though these biases

also depend on local time, AOT values, and surface properties. The results from this study

demonstrate that the assimilation of satellite aerosol retrievals not only improves the

²⁴ aerosol forecasts but also has the potential to reduce the uncertainties in modeling the

25 surface energy budget and other associated atmospheric processes. INDEX TERMS: 3337

26 Meteorology and Atmospheric Dynamics: Numerical modeling and data assimilation; 0305 Atmospheric

27 Composition and Structure: Aerosols and particles (0345, 4801); 3359 Meteorology and Atmospheric

28 Dynamics: Radiative processes; KEYWORDS: satellite assimilation in mesoscale model, GOES 8 aerosol

29 optical thickness, radiative effects

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33 1. Introduction

[2] Dust, a common aerosol over the desert, can be 34 transported to downwind areas thousands of kilometers 35 away from source regions [Reid et al., 2003], degrading 36 visibility and air quality, perturbing the radiative transfer in 37 the atmosphere [Hansen and Lacis, 1990], providing a 38 vector for disease-causing organisms [Shinn et al., 2000], 39 40 and exacerbating symptoms in people with asthma 41 [Prospero, 1999]. On the other hand, the atmospheric deposition of dust aerosols containing iron and other 42 trace elements is an important nutrient source for the 43 oceanic biota [Duce, 1991]. Both satellite remote sensing 44 [Christopher and Zhang, 2002; Wang and Christopher, 45 2003] and numerical models [Tegen and Fung, 1994; 46 Weaver et al., 2002] have been used to study dust radiative 47 effects and to monitor pollution. 48

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[3] Satellite remote sensing data sets are widely used to 49 map the geographical distribution of aerosols at high spatial 50 and temporal resolutions and to explore the effects 51 of atmospheric aerosols on the Earth's radiation budget 52 [Kaufman et al., 2002]. However, numerical models are 53 the preferred tool for studying the role of tropospheric 54 aerosols in modulating several important atmospheric 55 processes such as surface radiation energetics and 56 atmospheric heating rates [Carlson and Benjamin, 1980]. 57 Currently, satellite-derived aerosol information is not 58 commonly used in numerical models, especially regional 59 models. Using a case study of a dust event observed during 60 the Puerto Rico Dust Experiment (PRIDE) [Reid et al., 61 2003], we explore the utility of assimilating satellite-derived 62 aerosol information into numerical models to examine 63 aerosol radiative effects. 64

[4] The majority of Sun-synchronous polar-orbiting 65 satellites provides global images approximately once a 66 day and can only provide snapshots of large-scale aerosol 67 spatial distribution during the time of satellite overpass. On 68 the other hand, each geostationary satellite can view about 69

one third of the globe at temporal resolutions up to 15 min 70 or higher and is therefore suitable for the tracking, moni-71toring, and validation of regional-scale aerosol events 72[Wang et al., 2003]. Recent advances in satellite retrievals 73 can qualitatively capture aerosol optical thickness (AOT) 74distribution with high spatial [Kaufman et al., 2002] and 75 temporal resolutions [Christopher et al., 2002; Wang et al., 76 2003]. These satellite retrievals play an important role in the 77 estimation of aerosol radiative forcing [Christopher and 78 79 Zhang, 2002] and validation of numerical models [Chin et al., 2002]. However, current aerosol retrievals such as AOT 80 are mostly column quantities, thereby making it difficult for 81 examining the vertical distribution of aerosols and their 82 associated changes of radiative transfer in the atmosphere 83 [Kaufman et al., 2002]. 84

85 [5] Numerical models, both global and regional, are used 86 to forecast the three-dimensional (3-D) aerosol distributions. Compared to global simulations, regional mesoscale models 87 have finer spatial grids and therefore can be used in the 88 study of transport dynamics [Westphal et al., 1987, 1988; 89 Colarco et al., 2003], aerosol radiative forcing [Collins et 90 al., 2001], and modeling of air pollution [Binkowski and 91 Roselle, 2003]. Most regional models are off-line models 92[Byun and Ching, 1999], where the simulations are exter-93 nally driven by meteorological fields derived from other 94numerical models (such as Pennsylvania State University) 95National Center for Atmospheric Research Mesoscale Model 96 (NCAR MM5) or global climate models). Few regional-97 98 scale models explicitly include the radiative interactions in 99aerosol transport simulations, yet the radiative effect of dust aerosols can exert important effects on the forecast of 100 meteorological fields such as surface temperature and 101 boundary layer process, thereby motivating accurate char-102acterization of aerosols in model simulations [Carlson and 103Benjamin, 1980; Yu et al., 2002]. 104

[6] Limited availability of ground-based aerosol observa-105tions, especially over the oceans, constrains the accuracy of 106initial aerosol fields in both global and regional chemical 107transport models (CTMs) [Westphal and Toon, 1991; Chin 108 et al., 2002]. In general, in most numerical simulations, dust 109emission occurs when the wind speed over an erodible 110surface exceeds some threshold value. The choice of wind 111 speed threshold and parameterization of the emission flux is 112highly variable for different regions and is treated differ-113 ently in numerical modeling studies even for the same area 114 115[Takemura et al., 2000; Ginoux et al., 2001; Colarco et al., 2003]. Further challenges also exist in regional aerosol 116models that attempt to study the aerosol in one region but 117118 must span a large domain in order to capture the dust transport approaching from dust source region into the area 119of interest. This requires the use of nested grid configura-120121tions that add significant computational overhead to numerical simulations. 122

[7] Assimilation of satellite aerosol retrievals circumvents 123many of the problems associated with initializing aerosol 124 field in regional simulations. The assimilation of aerosol 125information derived from satellite measurements is a valu-126able tool to characterize the aerosol initial condition in the 127model, constrain the model simulation, and improve the 128model forecast [e.g., Collins et al., 2001; Yu et al., 2003]. 129130Previous studies [Collins et al., 2001; Yu et al., 2003] have 131assimilated the aerosol retrievals from polar-orbiting satel-



Figure 1. Model domain where the inset rectangle shows the domain of the fine grid. Also shown are the locations where the ground-based measurements were made at Roosevelt Road ((RR) 18.20° N, 65.60° W) and La Paguera ((LP) 17.97° N, 67.05° W).

lite into off-line models and showed that the assimilation of 132 satellite-retrieved AOT significantly improved the perfor- 133 mance of aerosol simulations. Although these studies uti- 134 lized satellite-derived AOT to represent aerosol spatial 135 distribution, they did not consider aerosol radiative effects 136 nor the possible atmosphere response to the aerosol radia- 137 tive effects such as change of atmospheric radiative heating 138 rate and surface radiative energy budget during the simula- 139 tion, since aerosol transport was modeled in an "off-line" 140 mode. Furthermore, polar-orbiting satellites used in these 141 studies have a repeat cycle of $1 \sim 2$ days which is insufficient 142 to capture the temporal evolution of the aerosol field. Field 143 experiments with intensive observations therefore provide a 144 good opportunity to investigate aerosol radiative effects in 145 numerical models. During PRIDE [Reid et al., 2003], half- 146 hourly dust AOT distribution over the ocean in the vicinity 147 of Puerto Rico was retrieved from the GOES 8 imager 148 [Wang et al., 2003]. Through assimilating GOES 8 AOT 149 into the Colorado State University Regional Atmospheric 150 Modeling System (CSU RAMS) [Pielke et al., 1992], and 151 by using observation data sets during PRIDE, this study 152 examines the radiative effects of dust aerosols. 153

2. Data and the Area of Study

[8] The area of study (Figure 1) is centered on Puerto 155 Rico, which was also the base for PRIDE [*Reid et al.*, 156 2003]. Of the five major dust events that were recorded 157 during PRIDE, we study the most severe event that occurred 158 during $20 \sim 23$ July 2000 (see Figure 2, A1–A5, and 159 detailed description in section 5.1). Profiles of aerosol 160 concentration from aircraft measurements, longwave and 161 shortwave downward radiative flux (W m⁻²) data from 162 Surface Measurements for Atmospheric Radiative Transfer 163 (SMART) [*Ji and Tsay*, 2000], Sun photometer inferred 164 AOT, and 2 m air temperature measurements made during 165 PRIDE are used. The location of two Sun photometers are 166

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Figure 2. Dust AOT retrieved from GOES 8 (A1–A5) and simulated AOT from the RAMS with (B1–B5) and without (C1–C5) assimilation of GOES 8 AOTs using nudging scheme. The percentage difference in AOT between aero-rad and noaero-rad cases is shown in D1–D5. The total downward radiative flux difference (longwave plus shortwave, W m^{-2}) at the surface and temperature difference (°C) in the model first layer above the surface are shown in E1–E5 and F1–F5, respectively. The black regions in A1–A5 are cloudy regions, and the white-outlined black areas are Puerto Rico land regions. See color version of this figure at back of this issue.

167 shown in Figure 1 including Roosevelt Road ((RR) 168 18.20° N, 65.60° W) and La Paguera ((LP) 17.97° N, 169 67.05° W). However, for this study, only AOT at LP is 170 used, since data at RR during $20 \sim 23$ July 2000 were not 171 available. Further details regarding the data sets are given 172 by *Reid et al* [2003].

173 [9] During PRIDE, dust AOTs retrieved from the 174 GOES 8 imager [*Wang et al.*, 2003] were used for 175 studying the diurnal change of dust forcing at the top 176 of atmosphere (TOA) and at the surface [*Christopher et* *al.*, 2003]. Using light scattering, absorption, and size 177 distribution measurements, the refractive index and single 178 scattering albedo of dust aerosols were estimated to be 179 1.53-0.0015i and 0.98 at 0.55 µm, respectively [*Wang et* 180 *al.*, 2003]. These aerosol optical properties are then used 181 in a discrete ordinate radiative transfer model to create 182 look-up tables for the GOES 8 AOT retrievals. The 183 satellite retrievals compared well with both in situ and 184 ground-based Sun photometer measurements [*Wang et al.*, 185 2003].

[10] The National Center for Environmental Prediction 187 (NCEP) reanalysis atmospheric data [Kalnay et al., 1996] at 188 0000, 0600, 1200 and 1800 UTC are used as a first guess 189for creating analysis of the meteorological fields used for 190specifying the initial conditions of the numerical model and 191for the evolution of the lateral boundary conditions for 192simulations starting at 1200 UTC on 20 July 2000 and 193ending at 1200 UTC on 23 July 2000. In addition, radio-194sonde and surface meteorological data obtained from data 195sets maintained at NCAR are also utilized in analysis of 196 these meteorological fields. Standard databases available 197 with the RAMS (version 4.3) are used to initialize 198topography and land use type at each grid cell while sea 199surface temperature is initialized using the National Oceanic 200and Atmospheric Administration (NOAA) satellite-derived 201202values (J. Vazquez et al., NOAA/NASA advanced very 203high resolution radiometer (AVHRR) Oceans Pathfinder Sea Surface Temperature Data Set User's Reference 204Manual Version 4.0, JPL Publication D-14070, available 205 at http://www.nodc.noaa.gov/woce V2/disk13/avhrr/docs/ 206usr_gde4_0_toc.htm). 207

3. Methodology 208

[11] The CSU RAMS (version 4.3) [Pielke et al., 1992] is 209modified to assimilate the GOES 8 AOT. Currently, the 210RAMS does not include dust aerosols in the radiative 211 transfer calculations. Therefore the RAMS was modified 212to include a sophisticated δ four stream radiative transfer 213model (84S RTM) [Liou et al., 1988; Fu and Liou, 1993] 214 that includes dust aerosol radiative properties measured 215during PRIDE [Christopher et al., 2003]. The aerosol 216transport model is built upon a tracer advection module in 217the RAMS combined with additional specification of source 218and sink mechanisms. 219

3.1. Description of the CSU RAMS and New 220

221**Modifications**

[12] The RAMS is a nonhydrostatic atmospheric model 222 [Pielke et al., 1992] that has been successfully used to 223simulate a wide range of atmospheric phenomenon includ-224 ing sea breezes, severe storms, flash flooding, downslope 225winds, air pollution, and atmospheric convection ranging 226227 from boundary layer cumulus to mesoscale convective systems [Cotton et al., 2002]. The RAMS utilizes finite 228difference approximations to solve conservation equations 229of mass, momentum, heat, and different phases of water on 230a polar stereographic grid in the horizontal and a terrain 231following sigma coordinate system in the vertical. At large 232 horizontal spatial scales, the RAMS uses convective 233parameterization schemes to account for precipitation 234mechanisms, while at smaller spatial scales it provides the 235capability to resolve cloud and precipitation processes 236through explicit bulk water parameterization. Surface layer 237parameterization along with a multilayer soil model and the 238sophisticated Land Ecosystem Atmosphere Feedback model 239(LEAF-2) [Walko et al., 2000] account for exchanges of 240241energy and momentum fluxes between the surface and the atmosphere. The RAMS also provides a wide range of 242243techniques for representing subgrid-scale turbulence and top and lateral boundary conditions. It includes a data 244analysis/assimilation module which blends upper air, 245

surface observations and gridded data from other models 246 such as NCEP reanalysis data to create products that are 247 used for initializing the model and nudging the top and 248 lateral boundaries of the model. The RAMS provides three 249 options of varying sophistication for longwave and short- 250 wave radiation calculations. However, none of three 251 existing radiation schemes in the RAMS accounts for the 252 radiative interactions of aerosols. Hence the 84S RTM is 253 implemented in the RAMS to compute the aerosol radiative 254 effects online. 255

3.1.1. The δ4S RTM

256

[13] The δ 4S RTM is a plane-parallel broadband radiative 257 transfer model, originally designed to calculate the radiative 258 flux at TOA and surface in clear and cloudy conditions [Fu 259 and Liou, 1993] and later modified for calculation of 260 radiative effect of aerosols, such as smoke [Christopher et 261 al., 2000] and dust [Christopher et al., 2003]. Gas absorp- 262 tion, water vapor absorption, and Rayleigh scattering are 263 included in the model calculations. The model divides the 264 shortwave (SW) spectrum (0.2 \sim 4 μ m) into six bands and 265 further divides the first band (0.2-0.7 µm) into 10 sub- 266 bands. The longwave (LW) spectrum is divided into 12 267 bands from 4 to 35.7 µm. To treat the radiative transfer 268 accurately, the δ 4S RTM uses a δ function to better represent 269 the phase function in the forward scattering direction [Liou 270 et al., 1988]. For the principal atmospheric gases the 271 difference between the δ 4S RTM and line-by-line irradiance 272 calculations is within 0.05% [Fu and Liou, 1993]. We have 273 modified the 84S RTM to include sea salt optical properties 274 [d'Almeida et al., 1991] (see section 3.1.3) and dust optical 275 properties derived from PRIDE [Christopher et al., 2003]. 276 Our recent studies indicate an excellent agreement between 277 calculated and observed downward shortwave irradiance at 278 the surface, with differences of <3% when aerosol effects 279 are carefully considered in the 84S RTM calculations 280 [Christopher et al., 2003]. 281282

3.1.2. Aerosol Transport

[14] The RAMS provides a generalized framework for 283 advection and diffusion of three-dimensional scalar quanti- 284 ties. In addition to scalar fields such as temperature and 285 water vapor as routinely used in standard computations, the 286 RAMS allows for specification of up to a hundred addi- 287 tional scalars. The aerosol processes, including emission, 288 advection, and deposition, can be expressed as 289

$$\frac{\partial C}{\partial t} = -u\frac{\partial C}{\partial x} - v\frac{\partial C}{\partial y} - w\frac{\partial C}{\partial \sigma} - \frac{\partial}{\partial x}\left(K_H\frac{\partial C}{\partial x}\right) - \frac{\partial}{\partial y}\left(K_H\frac{\partial C}{\partial y}\right) - \frac{\partial}{\partial \sigma}\left(K_L\frac{\partial C}{\partial \sigma}\right) + S,$$
(1)

where C is the aerosol concentration; u, v, and w are the 3-D 291 wind components; x, y, and σ denote 3-D coordinates; S 292 denotes net source/sink; K_L and K_H are the vertical and 293 horizontal diffusion exchange coefficient, and t is time. 294 While the advection module including diffusion exchange 295 of scalar variables is already available in the RAMS, new 296 emission/deposition parameterizations are incorporated into 297 the model for aerosol transport. 298

[15] Sea salt and dust aerosols are two primary types of 299 aerosols in the atmosphere during PRIDE [Reid et al., 300 301 2003]. In this study, the sea salt concentrations are 302 diagnosed as a function of wind speeds near the ocean 303 surface by using the following formula [*Blanchard and* 304 *Woodcock*, 1980; *Collins et al.*, 2001]:

$$C(z) = 5(6.3 \times 10^{-6} z)^{(0.21 - 0.29 \log_{10} U_{10})},$$
(2)

where C(z) is the sea salt concentration (in units of μ g m⁻³) at height z (in meters) above sea level and U_{10} is the 10 m wind speed (m s⁻¹). Above 300 m the profile decreases exponentially with a 500 m scale height.

[16] During the long-range transport of Saharan dust to 310 the Puerto Rico regions, large particles (diameter $>10 \ \mu m$) 311 are deposited either in the source regions or in the ocean 312 near the West African coast. Schütz and Jaenicke [1974] 313 314 found that nearly 75% of the large dust particles are deposited in the source area, leaving only 25% to reach 315the ocean 1500 km away. During PRIDE, Maring et al. 316 317 [2003] found that dust particles larger than 7.3 µm are mostly removed during the transport. Further analysis 318 319 showed that the vertical variations of the normalized dust size distribution are usually small during the dust events in 320 PRIDE [Reid et al., 2003], implying that it is reasonable to 321 simulate the mass concentration of dust aerosols in the 322 model rather than simulate the size distribution and use 323 the effective radius as the measure to calculate the gravita-324 325tional settling. This is especially useful for a case study in temporal scales of 2-3 days. The current model considers 326 327 the dry deposition process outlined by Slinn and Slinn [1980]. Wet deposition process is not included in the 328 simulation because no precipitation was recorded during 329 330 the period of study. The spatial distribution of GOES 8 AOT is used to initialize the 3-D aerosol concentration fields in 331 the RAMS. The procedure for deriving the 3-D aerosol 332 concentrations from GOES 8 AOTs is described in detail in 333 section 3.2. 334

335 3.1.3. Aerosol Optical Property Model

In [17] The aerosol optical property model is used to convert aerosol mass concentration into column AOT, thus providing a method for comparison between satellite and simulated AOT. In addition, the spatial distribution of AOT in each model layer is required in radiative transfer calculations. The relationship between mass concentration and AOT can be expressed as

$$\tau = \sum_{i=1}^{K} \left(\tau_{i,\text{dust}} + \tau_{i,\text{salt}} \right) = \sum_{i=1}^{K} \left(\mathcal{Q}_{\text{ext,dust}} C_{i,\text{dust}} + \mathcal{Q}_{\text{ext,salt}} C_{i,\text{salt}} \right) \Delta z_i,$$
(3)

where *i* is the index for vertical layers, *K* is the total number 344 of layers in the model, C is the mass concentration of 345 aerosols (g m⁻³), Q_{ext} is the mass extinction coefficient (m² g⁻¹), and Δz_i is the thickness (m) between different 346 347 layers. In this study, wavelength-dependent dust radiative 348 349 properties (e.g., single scattering albedo, asymmetric factor, extinction cross section, and mass extinction efficiency) in 350the shortwave spectrum (0~4 μ m) derived from PRIDE 351[Wang et al., 2003; Christopher et al., 2003] and in the 352 longwave spectrum from d'Almeida et al. [1991] are used. 353 For sea salt aerosols the complex refractive indices from 354

A. A. Lacis (Database of Aerosol Spectral Refractive 355 Indices, Global Aerosol Climatology Project, data available 356 at http://gacp.giss.nasa.gov/data_sets/), the sea salt size 357 distribution during PRIDE from *Maring et al.* [2003], and 358 the geometric growth factor at different relative humidity 359 due to the hygroscopic effect from *d'Almeida et al.* [1991] 360 are utilized in Mie calculations to compute the sea salt 361 radiative properties. 362

3.2. Assimilation of GOES 8 AOT

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[18] The model initialization of aerosol field at regional 364 scales is a challenging process, especially when the study 365 area is not in the aerosol source region. There are two ways 366 to tackle this problem. In the first method the aerosol 367 concentration from large-scale models can be used as initial 368 conditions for the mesoscale model [Colarco et al., 2003]. 369 However, compared to normal meteorological observations, 370 the current in situ aerosol observations are limited. There- 371 fore large uncertainties in dust concentrations could be 372 introduced when using output from global aerosol models 373 as initial conditions for mesoscale aerosol models for short- 374 term aerosol simulations. The second method is to use 375 aerosol information from satellite retrievals as the initial 376 condition [Westphal and Toon, 1991]. The advantage of this 377 second method is that the satellite data better represent the 378 aerosol spatial distributions, especially over large areas in 379 cloud-free conditions. The disadvantage is that current 380 satellite aerosol retrievals do not contain vertical informa- 381 tion and therefore it is difficult to infer the 3-D aerosol field 382 from 2-D satellite retrievals. As a consequence the shape of 383 aerosol vertical profiles together with a model of aerosol 384 optical properties (section 3.1.3) is required to convert the 385 2-D AOT into the 3-D aerosol mass concentrations. This 386 study uses the shape of the aerosol concentration profile as 387 measured by aircraft instruments [Reid et al., 2003]. The 388 initial dust profile in each model grid column has the same 389 shape as that derived from aircraft data but has different 390 mass concentrations such that when converting them 391 into AOT using equation (3), the AOT matches the 392 GOES 8 AOT at that grid point. Note, in equation (3), sea 393 salt concentration is diagnosed from wind fields, and the 394 only unknown variable is dust concentration. With known 395 vertical profile shape and GOES 8 AOT on the left side of 396 equation (3) the dust concentration profile for each grid 397 column can be computed. 398

[19] The justification for the above initialization method 399 is as follows. First of all, it is desirable to have the shape of 400 aerosol vertical profile in the conversions from 2-D AOT 401 field to the 3-D aerosol mass concentration field. The data 402 from aircraft measurements represent the most accurate 403 description of the aerosol vertical profile, especially for 404 the current study where the study area is relatively small and 405 dust vertical profile could have small spatial variations after 406 its long-distance transport across the Atlantic. Therefore the 407 aerosol profile from aircraft measurements is valuable and 408 cannot be ignored in the model initialization. Second, our 409 purpose in this study is to see how the direct consideration 410 of aerosol radiative effects in the model will affect the 411 simulation. In this context, column AOT is more important 412 than the vertical distribution itself. Studies have shown that 413 for aerosols in the lower boundary (<3 km) the vertical 414 distribution of aerosols has little effect on the calculation of 415

494

shortwave radiative energy at the surface as long as the
column AOT is the same [*Liao and Seinfeld*, 1998].
Therefore using aircraft vertical profile for the initialization
purposes is appropriate for this particular case.

420 [20] After the initialization the mass concentration of 421 aerosols is simulated, and AOT is predicted in each model 422 time step. To assimilate the GOES 8 AOT into the model 423 this study uses a Newtonian nudging scheme where 424 the values of model-predicted AOT tendency are adjusted 425 and the simulated AOTs are relaxed toward the satellite-426 retrieved AOT:

$$\left[\frac{\partial \tau'_{\text{mod-dust}}}{\partial t} \right]_{t} = (1 - \varepsilon) \left[\frac{\partial \tau_{\text{mod-dust}}}{\partial t} \right]_{t} + \varepsilon \frac{[\tau_{\text{GOES}}]_{t+1} - [\tau_{\text{mod-salt}} + \tau_{\text{mod-dust}}]_{t}}{\delta \Delta t}.$$
(4)

[21] In equation (4), $\partial \tau'_{\text{mod-dust}} / \partial t$ is the updated tendency of simulated dust AOT for the time step t after the 429430assimilation, $\tau_{\text{mod-dust}} / \partial t$ is the tendency of dust AOT at 431 time step *t* before the assimilation, and $\tau_{\text{mod-salt}}$ is the optical 432thickness of sea salt. The GOES 8 AOT at time step t + 1433 (e.g., $[\tau_{GOES}]_{t+1}$) is computed by linear interpolation of 434 GOES 8 AOT from the two assimilation time periods (see 435section 4). The difference between $[\tau_{GOES}]_{t+1}$ and the 436 simulated AOT (sum of $\tau_{mod-salt}$ and $\tau_{mod-dust})$ at time step 437t divided by the integration time interval Δt is the AOT time 438tendency if we want to force the modeled AOT to match the 439GOES 8 AOT in the next step t + 1 (recall $\tau_{t+1} = \tau_t +$ 440441 tendency $\times \Delta t$). This tendency is then adjusted by a 442confidence factor (δ), where smaller δ (nearly equal to 1) implies that GOES 8 AOT has very high accuracy, hence 443larger confidence of this tendency term. The ε value in 444 equation (4) is a 2-D weighting factor that is used in the 445RAMS to control the effect of the nudging term as a 446 function of spatial location. Notice, for $\delta = 1$ and $\varepsilon = 1$ 447 the effect of the modeled AOT tendency $\partial \tau_{mod-dust} / \partial t$ will 448 be neglected, and equation (4) forces the modeled AOT to 449exactly relax toward the GOES 8 AOT. We set $\delta = 1.02$ 450because GOES 8 AOT has a slight overestimate bias when 451compared to ground and in situ measurements [Wang et al., 4522003]. Since this study does not have aerosol fields outside 453the GOES 8 AOT coverage domain, it is desirable for the 454nudging term to have maximum effect along the lateral 455456boundaries (to provide the aerosol source) while letting the 457effect of nudging be minimized in the model interior, allowing for the aerosol transport parameterization. Such 458 behavior is obtained by specifying ε to be nearly one for the 459five nearest grid points to the lateral boundaries while 460exponentially decreasing it to zero at grid points in the 461model interior. 462

[22] Through the assimilation of GOES 8 AOT the 463modeled AOT tendency is then optimally modified using 464equation (4), and consequently, the modeled dust concen-465 tration tendency in each layer needs to be adjusted in order 466 to be consistent with the updated AOT tendency. Since the 467AOT is a column quantity, assumptions on the aerosol 468 vertical distribution must be made in order to adjust the 469 tendency of mass concentration in each layer. In the 470assimilation of advanced very high resolution radiometer 471(AVHRR) AOT in CTMs, Collins et al. [2001] assumed that 472

the shape of aerosol vertical distribution is the same before 473 and after the assimilation, and thus a scale factor η can be 474 obtained for each model layer *i*: 475

$$\eta = \left[\frac{\tau'_{\text{mod-dust}}}{\tau_{\text{mod-dust}}}\right]_{t+1} = \left[\frac{C'_{i,\text{dust}}}{C_{i,\text{dust}}}\right]_{t+1},\tag{5}$$

where $C'_{i,\text{dust}}$ and $C_{i,\text{dust}}$ are the mass concentration of dust 477 after the assimilation and before the assimilation at layer *i*, 478 respectively. Therefore from equation (5) the new dust 479 concentration tendency after the assimilation can be 480 calculated using 481

$$\begin{bmatrix} C_{i,\text{dust}} \end{bmatrix}_{t} + \begin{bmatrix} \frac{\partial C'_{i,\text{dust}}}{\partial t} \end{bmatrix}_{t} \Delta t = \eta \begin{bmatrix} (C_{i,\text{dust}})_{t} + \left(\frac{\partial C_{i,\text{dust}}}{\partial t}\right)_{t} \Delta t \end{bmatrix}$$

$$\begin{bmatrix} \frac{\partial C'_{i,\text{dust}}}{\partial t} \end{bmatrix}_{t} = \eta \begin{bmatrix} \frac{\partial C_{i,\text{dust}}}{\partial t} \end{bmatrix}_{t} + \frac{\eta - 1}{\Delta t} \begin{bmatrix} C_{i,\text{dust}} \end{bmatrix}_{t}.$$
(6)

[23] In summary, aerosol initial fields are obtained by 484 converting the 2-D GOES 8 AOT into 3-D fields with the 485 assumption that the vertical profiles are the same as 486 prescribed by aircraft measurements in all grids. In the 487 following simulations the dust concentration field is 488 adjusted so that the modeled AOT is relaxed toward the 489 GOES 8 AOT. Therefore except in the initial time step the 490 vertical profiles of aerosol concentration change with time 491 in the different grids. 492

4. Model Configuration and Initialization

[24] The numerical simulation in this study utilizes a 495 nested grid configuration: a fine grid of 34×34 points 496 and 40 km grid spacing covering the island of Puerto Rico, 497 nested within a coarse grid with 32×32 grid points and 49880 km grid spacing (Figure 1). Both the configurations use a 499 stretched vertical grid of 30 grid points and a grid stretch 500 ratio of 1.2, with the vertical grid spacing increasing from 501 100 m near the surface to a maximum of 750 m higher up in 502 the atmosphere. All numerical simulations used in this study 503 are initiated at 1200 UTC on 20 July 2000 and end at 504 1200 UTC on 23 July 2000. Time steps of 60 and 10 s are 505 used for time stepping the coarse and fine grids, 506 respectively. Analysis of meteorological fields derived from 507 NCEP reanalysis gridded pressure level data [Kalnay et al., 508 1996] as well as upper atmospheric and surface observa- 509 tions at 1200 UTC on 20 July 2000 is used to initialize 510 atmospheric conditions in the model simulation. Sea surface 511 temperature is initialized using the AVHRR sea surface 512 temperature (SST) product (J. Vazquez et al., NOAA/NASA 513 advanced very high resolution radiometer (AVHRR) Oceans 514 Pathfinder Sea Surface Temperature Data Set User's 515 Reference Manual Version 4.0, JPL Publication D-14070, 516 available at http://www.nodc.noaa.gov/woce V2/disk13/ 517 avhrr/docs/usr gde4 0 toc.htm) available at a temporal 518 frequency of 7 days. The SST is assumed to vary linearly 519 with time for the period in between the observations. 520 Similar meteorological analyses at 6 hour intervals are 521 used to nudge the model's lateral and top boundaries. The 522 GOES 8 AOT is produced at 4×4 km [Wang et al., 2003]. 523 Three-dimensional aerosol concentration fields derived 524

from GOES 8 data valid at 1200 UTC are then averaged 525into 40 and 80 km grids to initialize aerosol fields in the 526model. Although GOES 8 AOT has a temporal resolution of 527a half hour, the model only assimilates the GOES 8 AOT 528twice per day, one at 1331 UTC and one at 1931 UTC 529(on the first day, only 1931 UTC is assimilated). There are 530several reasons for this implementation. First, the GOES 8 531AOT at certain time periods are needed for intercomparison 532with model simulations; otherwise it is difficult to judge the 533performance of a model only on the basis of the comparison 534with ground-based point measurements [Chin et al., 2002]. 535Second, if the model assimilates AOT every half hour, the 536model simulation will lack its value and will be forced too 537much toward the GOES 8 AOT. Finally, our near-future 538goal is to assimilate the MODIS AOT from both Terra and 539540Aqua [Kaufman et al., 2002] into the model, and assimila-541tion of GOES 8 AOT twice per day provides a prototype for our future studies. 542

[25] Two different numerical simulations are considered 543in this study, and they differ only in the treatment of aerosol 544radiative effects, i.e., inclusion of aerosol effects (aero-rad 545case) and no inclusion of aerosol effects (noaero-rad case). 546In addition, a "free-run" simulation without any assimila-547tion of GOES 8 AOT but with direct online consideration of 548aerosol radiative effects is also performed to test the impact 549of assimilation. 550

551 5. Results

552[26] The AOT from the numerical simulations is compared against the satellite-derived AOT and point measure-553ments of AOT derived from ground-based Sun photometer. 554Surface measurements of downward shortwave and long-555wave fluxes and 2 m air temperature are compared against 556the simulated values. The differences of AOT and surface 557 radiative energy budget between aero-rad and noaero-rad 558cases are further compared and analyzed over the whole 559simulation domain at different time periods. 560

561 5.1. Comparison of Modeled AOTs

[27] Satellite observations show that dust approached 562 Puerto Rico during the morning of 20 July 2000 (Figure 2, 563A1) and reached the island of Puerto Rico in the evening 564(Figure 2, A2). The dust layer then passed through the 565island between late evening on 20 July 2000 (Figure 2, A3) 566to late evening on the second day (Figure 2, A4) and ended 567 by 22 July 2000 (Figure 2, A5). The modeled AOT without 568569any assimilation are shown in Figure 2, B1-B5. In the absence of nudging along the lateral boundaries, the model 570does not account for external transport of dust into the 571572computational domain. In this case, as shown in Figure 2, B1-B5, the dust layer moves quickly across the model 573domain and disappears completely in one day. This is a 574 typical behavior for a limited area model without proper 575specification of boundary conditions. 576

577 [28] With assimilation the two numerical model simula-578 tions (aero-rad and noaero-rad cases) exhibit spatial patterns 579 of AOT similar to GOES 8 observations (only aero-rad case 580 is shown in Figure 2, C1-C5). Simulated AOT fields from 581 the second grid of the aero-rad case show a very similar 582 sequence of events (Figure 2, C1-C5), and the location of 583 dust front in the model simulation agrees well with that from



Figure 3. Simulated AOT versus Sun photometer AOT (dots) at La Paguera. Also shown is the GOES 8 AOTs (squares). Vertical dotted line shows the times (1331 and 1991 UTC on each day) when GOES AOT is assimilated into the model.

satellite-retrieved AOT. Note that the satellite-retrieved 584 AOT is plotted at a spatial resolution of 4 km, while the 585 simulated AOT field has a grid spacing of 40 km. Because 586 of these differences the simulated AOT field is smoother 587 compared to observations and lacks some of the observed 588 smaller-scale details (e.g., Figure 2, A3). 589

[29] Comparison of point observations of AOT derived 590 from Sun photometer measurement at LP against AOT from 591 the closest grid point in model simulations (Figure 3) shows 592 general agreement. The modeled AOT matches the Sun 593 photometer AOT and captures the temporal evolution of 594 dust event very well. Note that the model simulated values 595 are an average over a 40×40 km area while the observa- 596 tions are essentially point measurement that resolves fine- 597 scale features within the dust event. Also notice that the 598 GOES 8 AOT used in this study slightly overestimates the 599 Sun photometer AOT that can result in a positive bias in 600 modeled AOT (Figure 3). During the simulations a total of 601 seven GOES 8 AOT retrievals are assimilated, at 1331 and 602 1931 UTC as indicated by dotted vertical lines in Figure 3. 603 Though simple linear nudging technique (see equation (3)) 604 is used in this study, the nudging provides a correction for 605 the dynamical simulations in the models. Thus the final 606 results combine the strengths of both nudging corrections 607 and model simulations and therefore are not a simple linear 608 process. Such nonlinear features are distinct, as the modeled 609 AOT can capture the diurnal variations of AOT very well 610 (Figure 3), especially when the dust reaches Puerto Rico. 611 We argue that using linear nudging alone will not produce 612 such a feature since the modeled AOT is the composite 613 effect from both dynamical modeling and the correction 614 from satellite measurements. The implication of Figure 3 is 615 that the best estimation of AOT, especially the diurnal 616 change of AOT, should come from combined satellite 617 measurements and numerical simulations. 618

[30] Figure 3 also shows that the difference of modeled 619 AOT in aero-rad and noaero-rad cases is visually small, 620 though the AOT difference of ~ 0.03 does exist at some 621 time periods. The aerosol radiative effects on the simulation 622



Figure 4. Comparison between measured and modeled downward fluxes at Roosevelt Road in aero-rad and noaero-rad cases for (a) shortwave (SW) and (c) longwave (LW) flux. The differences of modeled and measured downward fluxes are shown for (b) shortwave and (d) longwave fluxes. (e) Total flux difference (e.g., shortwave difference plus longwave difference) between the two simulation cases is shown. All circles denote the in situ data measured from the Surface Measurements for Atmospheric Radiative Transfer (SMART) instrument suite. (f) Comparison between modeled 2 m air temperature with the measured values is shown. See color version of this figure at back of this issue.

of AOTs are then further investigated in terms of the relative 623 difference (in percentage) of AOT in two simulations (e.g., 624 $(\tau_{noaero-rad} - \tau_{aero-rad})/\tau_{aero-rad} \times 100\%$, Figure 2, D1–D5) 625in the whole model domain. Overall, the relative difference 626 is within $-5 \sim 10\%$, and the lack of consideration of aerosol 627 radiative effect during the simulation tends to produce a 628 positive bias of AOT. The largest differences occur in areas 629 with low AOTs dominated by sea salt. The maximum 630 absolute difference of AOT is ~ 0.05 in areas where dust 631 is dominant. As will be shown later, the noaero-rad case 632 overestimates the total downward flux at the surface, and 633 such overestimation causes different atmosphere responses 634 which then leads to differences in modeled AOTs. Even 635 though the atmosphere system responds to aerosol forcing 636 in multiple ways, this study focuses on the impact of dust 637 638 aerosols on surface radiative energy budget.

5.2. Downward Flux Comparisons 639

[31] Figure 4a shows the comparison of modeled and 640 measured downward solar flux at the surface at Roosevelt 641 Road, and their difference is shown in Figure 4b. Though 642 643 both simulations tend to overestimate the downward fluxes. the overestimation in the aero-rad case is much smaller than 644 645that in the noaero-rad cases. Note that the measured SW flux only covers a section of the solar spectrum from 646 $0.28 \sim 2.8 \ \mu m$, while the modeled flux shown in Figure 4b 647 covers the whole SW spectrum (e.g., $0 \sim 4 \mu m$). This 648 partially explains the difference of observed flux and 649 modeled fluxes in the aero-rad case. Using the same dust 650 optical properties and Sun photometer AOTs, off-line 651calculations indicate that the difference between 84S RTM 652modeled SW flux (at 0.28~2.8 µm) and measured flux is 653 within 15 W m⁻² [*Christopher et al.*, 2003]. Uncertainties 654in the modeled AOTs are another factor that potentially 655656 contributes to overestimation of downward solar flux by the aero-rad case. Accounting for these two factors and other 657 658 factors such as the possible presence of cloud (e.g., a sudden 659 sharp drop of measured SW flux in the late afternoon of 22 July), the aero-rad case better represents the averaged 660 downward flux. Compared to the aero-rad simulation 661 and observations, the neglect of scattering by aerosols in 662 noaero-rad case leads to overestimation of instantaneous flux from 10 up to 100 W m⁻² depending on the magnitude 663 664 of dust AOTs and time of day (e.g., solar zenith angle, 665 Figure 4b). This is a common feature for most mesoscale 666 models that do not include radiative interactions of aerosols 667 [e.g., Chen and Dudhia, 2001]. Overall, the lack of consid-668 669 eration of aerosol radiative effect results in enhancement of solar energy at the surface, with an average daytime 670 "warming" bias of 40 W m⁻² that is also consistent with 671 672 the previous studies [Christopher et al., 2003].

673 [32] Dust absorbs in the longwave part of the electromagnetic spectrum [d'Almeida et al., 1991], and this effect is 674 obvious when comparing the model-simulated downward 675 longwave flux and surface observations (Figures 4c and 4d). 676 The dust layer absorbs the outgoing longwave flux from the 677 surface and reemits it back to the surface, thus increasing 678 the downward longwave flux. Both simulations underesti-679 mate the downward longwave fluxes (Figure 4c), but the 680 aero-rad case shows the least deviation from the observa-681 tions. Note that the simulated downward longwave flux 682 represents an area averaged over 40×40 km while 683

observations are essentially point measurement. The 684 uncertainties in the vertical distribution and dust properties 685 as well as the surface heterogeneity (see section 5.3) could 686 also contribute to the relatively large bias of modeled 687 longwave flux. Compared to the aero-rad simulation, the 688 noaero-rad simulations exhibit a "cool" bias (less down- 689 ward longwave) at the surface (Figure 4d). The difference is 690 about -10 W m⁻² depending on the magnitude of AOTs. 691

[33] The comparison of the total flux difference (short- 692 wave plus longwave) is highly variable depending on the 693 magnitude of AOT and the local time (Figure 4e). Figure 2, 694 E1-E5, shows that the total downward flux difference is 695 highly consistent with the transport of dust aerosols. 696 During the daytime the lack of aerosol radiation effect leads 697 to an overestimation of surface incoming energy from 698 $40 \sim 60 \text{ W m}^{-2}$, though the instantaneous values depend on 699 locations and local time. During the nighttime, however, the 700 surface incoming energy is underestimated $\sim 10 \text{ W m}^{-2}$ if 701 dust aerosol radiative effects are not considered. Since these 702 energy fluxes are crucial inputs for the land surface param- 703 eterization, the performance of radiative transfer parameter- 704 izations used in mesoscale models needs to be further 705 evaluated, which is beyond the scope of the current study. 706

5.3. Comparison of Air Temperatures

707 [34] The comparison of the modeled and measured 2 m 708 air temperature is shown in Figure 4f. Figure 4f shows a 709 consistent pattern where the model-simulated nocturnal 710 temperatures are lower than measurements. This is consis-711 tent with surface radiation energy budget analysis discussed 712 in section 5.2, which shows underestimation of downward 713 longwave flux, a dominant control on the nocturnal 714 evolution of surface air temperature. During daytime, 715 modeled temperature in both cases has a similar magnitude 716 wit the observations. Overall, since downward flux is 717 overestimated in noaero-rad case during daytime and 718 underestimated during nighttime, the temperature in 719 noaero-rad simulations is slightly higher than that of 720 aero-rad cases in daytime and lower at nighttime. However, 721 such differences are very small, and it is difficult to judge 722 which one is better if solely based on Figure 4e. The reason 723 for this small difference is due to the surface heterogeneity 724 associated with the model grid point covering the Roosevelt 725 Road location. Roosevelt Road is located on the eastern 726 edge of Puerto Rico (Figure 1), and the nearest grid point 727 covers a 40×40 km area that includes both land and ocean. 728 Compared to water bodies, the temperature in the boundary 729 layer over land is significantly more sensitive to changes in 730 surface radiative energy budget. Water bodies absorb down-731 ward radiative energy over a deep layer, and mixing 732 transports the energy to further depths. In addition, the high 733 heat capacity of water causes a comparatively slow change 734 in water temperature. In the RAMS, SST is not predicted 735 but is specified using the AVHRR ocean surface tempera-736 ture product (J.Vazquez et al., NOAA/NASA advanced very 737 high resolution radiometer (AVHRR) Oceans Pathfinder Sea 738 Surface Temperature Data Set User's Reference Manual 739 Version 4.0, JPL Publication D-14070, available at http:// 740 www.nodc.noaa.gov/woce V2/disk13/avhrr/docs/ 741 usr gde4 0 toc.htm). The RAMS uses linear interpolation 742 to account for variations in SST over a timescale of a week, 743 but the diurnal variations are not explicitly simulated. 744

However, note that the diurnal and short-timescale 745 $(2\sim3 \text{ days})$ variations in SST are expected to be negligible 746 because of reasons discussed above. Over land, the down-747748 ward radiation is absorbed by a thin layer of surface soil, and small soil heat capacity causes rapid changes in soil 749 temperature compared to a water surface. Convective 750mixing efficiently transports a major part of the energy 751 absorbed by the land surface to the boundary layer, resulting 752in it being more responsive to changes in surface radiative 753 energy budgets. Therefore at the grid point used for com-754 parison against measurements from Roosevelt Road, the 755 impact of aerosol radiative forcing on the 2 m air temper-756 ature is diminished because of the presence of ocean in that 757 grid cell. 758

[35] The aerosol radiative forcing effect on the boundary 759 760 layer temperature and its modulation by the nature of 761 surface type become more obvious when the temperature difference between two cases in the first model layer above 762 the ground is compared. Figure 2, F1-F5, shows that 763 such temperature difference patterns are correlated to 764corresponding spatial distribution patterns of AOT. Over 765 ocean, where changes in surface radiation energy budget 766 have negligible effect on boundary layer air temperature, 767 direct heating of air by the absorption of dust aerosols is the 768 dominant process. Therefore presence of atmospheric dust 769over ocean leads to direct warming of air, causing air 770 temperatures to be warmer in aero-rad case over such region 771 [Carlson and Benjamin, 1980]. However, over land, heat 772 773 transfer from the surface to the atmosphere (either through 774the sensible heating or the vertical turbulence convection) has the more dominant effect and overshadows the direct 775 radiative heating effect of the atmospheric dust. Decrease in 776 downward solar radiation and associated reduction in sen-777 sible heat transfer to the atmosphere over land in the 778 presence of atmospheric dust leads to air temperature being 779 780 lower in aero-rad case. The neglect of aerosol radiative effect leads to the temperature change in the first model 781 layer from -0.5° C over the ocean to 0.5° C over the land 782 when dust is dominant. 783

785 6. Discussion and Conclusion

[36] In this study, a method for assimilating the satellite 786 derived AOT into the regional mesoscale models is devel-787 oped. A four stream radiative transfer parameterization was 788 added into the RAMS to explicitly consider the aerosol 789 radiative effects during the simulation. Through the com-790 791 parison with in situ, ground-based and satellite observations it is found that the inclusion of aerosol radiative effects 792 793 improves the overall performance of the modeled aerosol 794 fields and surface radiation budgets, though improvement of 795 2 m air temperature is minimal because of the relatively coarse grid size that cannot resolve the detailed surface 796 characteristics near the observation site. 797

[37] The implication of this study is twofold. First, for 798 moderate to high aerosol loadings, aerosol radiative effects 799and atmospheric response to these effects are significant 800 enough to be considered in the simulation of aerosol 801 transport and weather forecast. Off-line simulations without 802 proper treatment of aerosol radiative feedbacks may exhibit 803 biases which could be severe depending on the aerosol type 804 and loading as well as surface type, consistent with a recent 805

theoretical study which shows that aerosol absorption in the 806 atmosphere could alter the surface radiation energy budget 807 and the profile of heating rate significantly enough to 808 influence the vertical diffusions in the boundary layer [Yu 809 et al., 2002]. Second, since there are only a few observa- 810 tions available that routinely measure the aerosol concen- 811 tration over the globe, assimilation of satellite aerosol 812 retrievals into the mesoscale numerical models provides 813 an important tool to narrow the uncertainties in aerosol 814 source function and has the potential to become a cost- 815 effective method to improve particulate matter forecast, 816 especially in places where ground-based observations are 817 sparse. One of the major obstacles in this type of assimila-818 tion is the lack of information on the vertical structures of 819 aerosol distributions in current satellite aerosol retrievals. 820 Aerosol vertical profiles from spaceborne lidar measure- 821 ments [Winker et al., 2002] could provide valuable infor- 822 mation for the assimilation of satellite-derived AOTs in the 823 near future. Therefore with direct consideration of aerosol 824 radiation effects and assimilation of satellite aerosol retriev- 825 als the aerosol transport and distribution can be more 826 realistically simulated, which also has the potential to bring 827 overall improvement to weather forecast. 828

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Figure 2. Dust AOT retrieved from GOES 8 (A1–A5) and simulated AOT from the RAMS with (B1–B5) and without (C1–C5) assimilation of GOES 8 AOTs using nudging scheme. The percentage difference in AOT between aero-rad and noaero-rad cases is shown in D1–D5. The total downward radiative flux difference (longwave plus shortwave, W m^{-2}) at the surface and temperature difference (°C) in the model first layer above the surface are shown in E1–E5 and F1–F5, respectively. The black regions in A1–A5 are cloudy regions, and the white-outlined black areas are Puerto Rico land regions.



Figure 4. Comparison between measured and modeled downward fluxes at Roosevelt Road in aero-rad and noaero-rad cases for (a) shortwave (SW) and (c) longwave (LW) flux. The differences of modeled and measured downward fluxes are shown for (b) shortwave and (d) longwave fluxes. (e) Total flux difference (e.g., shortwave difference plus longwave difference) between the two simulation cases is shown. All circles denote the in situ data measured from the Surface Measurements for Atmospheric Radiative Transfer (SMART) instrument suite. (f) Comparison between modeled 2 m air temperature with the measured values is shown.